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## Single-crystal growth and superconducting state of $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$

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### Abstract

Single crystals of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  were grown by a CsCl-flux method. Electrical resistivity  $\rho(T)$  measurements were performed to reveal its superconducting properties. The  $\rho(T)$  of the single crystal shows semiconducting-like behavior and superconducting transition at 3.4 K. The value is slightly higher than that of the polycrystalline sample with substitution amounts of  $x = 0.2$  ( $T_c \sim 2.5$  K). From  $\rho(T)$  measurements in several magnetic fields,  $\mu_0 H_{c2}^{\parallel ab}(0)$  and  $\mu_0 H_{c2}^{\parallel c}(0)$  are estimated to be 17.2 T and 0.59 T, respectively. The superconducting anisotropic parameter  $\gamma$  is determined to be 29.2.

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**Keywords:** superconductor; BiS<sub>2</sub>-based compound

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### 1. Introduction

Since the discovery of superconductivity with superconducting transition temperature  $T_c$  at 8.6 K in the layered bismuth oxysulfide  $\text{Bi}_4\text{O}_3\text{S}_3$  [1], other BiS<sub>2</sub>-based superconductivity have been studied.  $\text{LaOBiS}_2$  crystallizes with a space group  $P4/nmm$  (No. 129) and this structure is composed of alternating superconducting BiS<sub>2</sub> layers and blocking LaO layers, which are similar to Fe-based superconductors. Substituting F for O induces superconductivity with  $T_c = 2.5$  K [2], as well as tetravalent elements (Ti, Zr, Hf and Th) for La [3]. In addition, by replacing S atoms

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to Se atoms,  $T_c$  is enhanced:  $T_c = 4.2$  K for  $\text{LaO}_{0.5}\text{F}_{0.5}\text{BiS}_2$  [4], 3.7 K for  $\text{LaO}_{0.5}\text{F}_{0.5}\text{BiSe}_2$  [5]. The enhancement of  $T_c$  in Se doped compounds is induced by an in-plane chemical pressure [6].

In this article, we report the single crystal growth of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  and its superconducting properties. In order to clarify superconducting state, we measured the electrical transport properties of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  at low temperature and strong magnetic fields up to 11 T.

## 2. Experimental Details

Single crystals of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  were grown by a CsCl-flux method similar to the previous reports [5,6]. The starting materials of 0.8 g and CsCl of 5 g were mixed and sealed in an evacuated quartz tube. The tube was heated up to 880 °C in 10 h and to 900 °C in 2 h, then kept for 10 h and cooled down to 650 °C for 125 h. The obtained materials were washed by water, ethanol, and acetone in order to remove the flux materials. The obtained single crystals were plate-like shape with approximately  $1.0 \times 1.0 \text{ mm}^2$ . The crystal structure of the single crystal was examined by an X-ray diffraction (XRD) method using a conventional X-ray spectrometer equipped with  $\text{Cu-K}\alpha$  radiation and a monochromator (RAD-2X, Rigaku). Electrical resistivity  $\rho(T)$  from 0.5 K to 300 K was measured under magnetic fields up to 11 T along the *ab*-plane and *c*-plane by dc-four-probe method in a  $^3\text{He}$  cryostat (Oxford Heliox VL).

## 3. Experimental Result

### 3.1. X-ray diffraction

Figure 1 shows the XRD diffraction pattern of several single crystals  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  at room temperature. Only (00*l*) diffraction peaks were observed, indicating the crystallographic *c*-axis is perpendicular to the crystal plane. All the reflections can be indexed as the space group of *P4/nmm* and no extra peaks due to impurity phase can be detected from the XRD pattern. To estimate lattice constants, powder X-ray diffraction was performed by using powder crashed from single crystals. The lattice constants of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  are obtained to be  $a = 4.0856 \text{ \AA}$  and  $c = 13.453 \text{ \AA}$ . These values are slightly larger than those of the single crystal  $\text{LaO}_{0.5}\text{F}_{0.5}\text{BiS}_2$ , which suggest that S atoms are substituted with Se atoms.

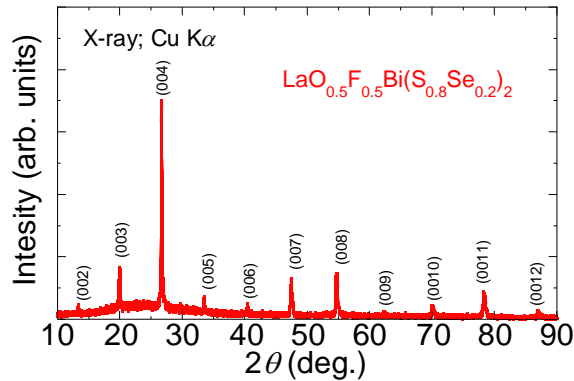


Fig. 1. XRD diffraction pattern of single crystals  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  at room temperature.

### 3-2. Electrical resistivity

Figure 2 shows temperature dependence of the electrical resistivity  $\rho(T)$  of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  down to 0.5 K with electrical current along the *ab*-plane. A semiconducting-like behavior was observed and superconducting transition was detected.  $T_c$  is defined as 50 % drop from residual resistivity  $\rho_0$  and determined to be 3.4 K.  $\rho_0$  is defined at just above  $T_c$  ( $\rho_0 = 3.5 \text{ m}\Omega\text{cm}$ ). Compared with polycrystalline sample of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  [7],

although the behavior of  $\rho(T)$  and the  $\rho_0$  is almost the same value,  $T_c$  is slightly higher than that of polycrystalline sample ( $T_c \sim 2.5$  K).

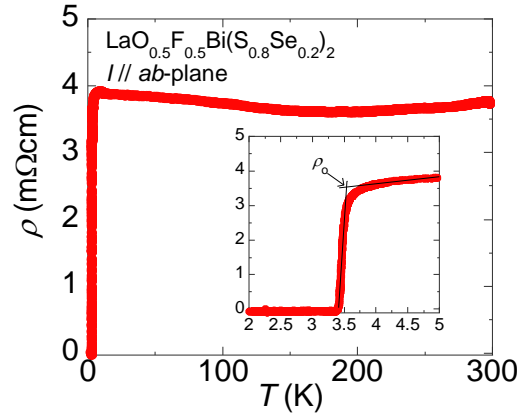


Fig. 2. Temperature dependence of the electrical resistivity  $\rho(T)$  of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$ . The Inset shows  $\rho(T)$  near  $T_c$ .

Figure 3 shows temperature dependence of electrical resistivity  $\rho(T)$  of single crystal  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  in several magnetic fields along  $ab$ -plane (a) and  $c$ -plane (b).  $T_c$  shifts toward lower temperature with an increase of the magnetic field. Superconductivity is robust against magnetic field applying along  $ab$ -plane whereas it is quite sensitive to the magnetic field applying along  $c$ -plane.

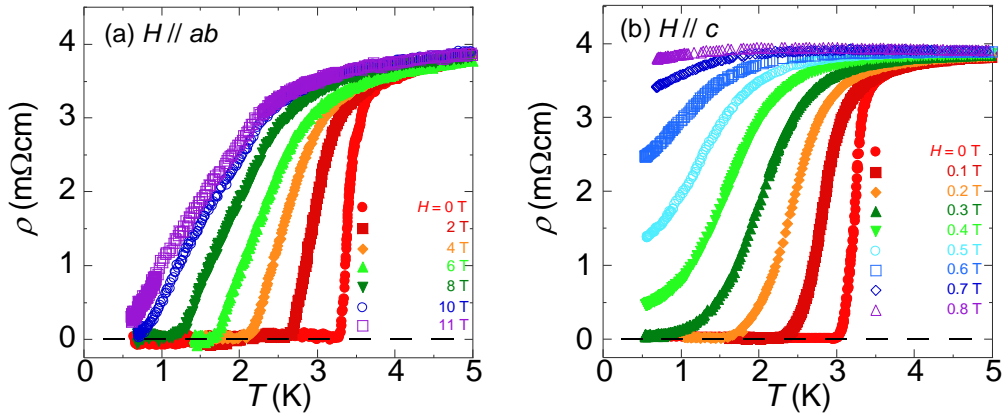


Fig. 3. Temperature dependence of the electrical resistivity  $\rho(T)$  of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  in several magnetic fields along  $ab$ -plane (a) and  $c$ -plane (b) at low temperatures.

Figures 4(a) and (b) show temperature dependence of upper critical field  $\mu_0 H_{c2}^{\parallel ab}(T)$  and  $\mu_0 H_{c2}^{\parallel c}(T)$ . Error bars are determined by the temperature of 10 % and 90 % drop from  $\rho_0$ . For  $ab$ -plane, the  $\mu_0 H_{c2}^{\parallel ab}(T)$  curve shows non- $T$ -linear behavior. A Werthamer-Helfand-Hohenberg (WHH) fitting in a dirty limit, as presented by dashed line, is not completely correspond to experimental data. The type-II SC approximation  $\mu_0 H_{c2}(T) = \mu_0 H_{c2}(0)(1-(T/T_c)^{3/2})^{3/2}$  [8], presented by solid line, is in good agreement with the data.  $\mu_0 H_{c2}^{\parallel ab}(0)$  is estimated to be 17.2 T. For  $c$ -plane, the  $\mu_0 H_{c2}^{\parallel c}(T)$  curve shows nearly  $T$ -linear behavior. Because the WHH fitting and type-II SC approximation are not corresponding to the data, we fitted  $\mu_0 H_{c2}^{\parallel c}(T)$  using polynomial equation  $\mu_0 H_{c2}(T) = \sum_{n=0}^3 a_n T^n$  presented by solid line.  $\mu_0 H_{c2}^{\parallel c}(0)$  is estimated to be 0.59 T. The superconducting anisotropic parameter  $\gamma = H_{c2}^{\parallel ab}(0)/H_{c2}^{\parallel c}(0)$  is determined to be 29.

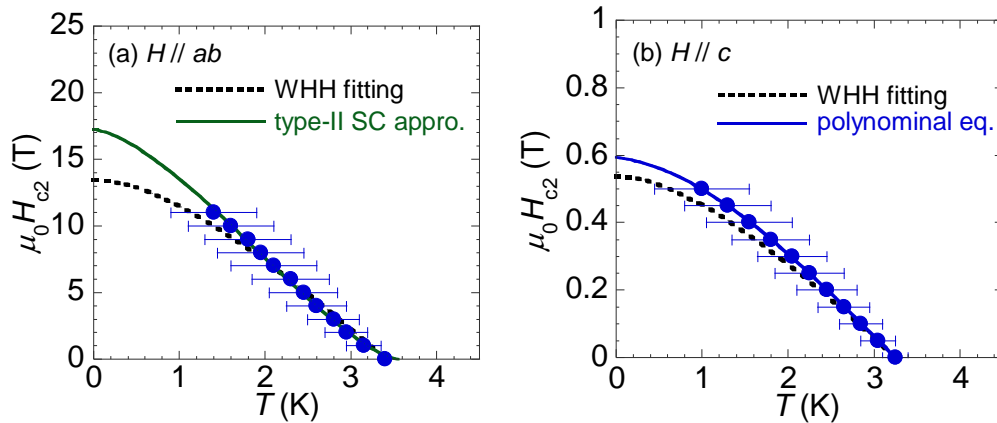


Fig. 4(a) Temperature dependence of the upper critical field  $\mu_0 H_{c2}(T)$  of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  in several magnetic fields along  $ab$ -plane. The dotted line shows the WHH prediction in a dirty limit. The Solid line represents type-II SC approximation. (b). Temperature dependence of  $\mu_0 H_{c2}(T)$  along  $c$ -plane. The dotted line shows the WHH prediction in a dirty limit. Solid line represents polynomial equation.

#### 4. Conclusion

We succeeded in growing the single crystal of  $\text{LaO}_{0.5}\text{F}_{0.5}\text{Bi}(\text{S}_{0.8}\text{Se}_{0.2})_2$  by a CsCl-flux method. The single crystal sample exhibits superconducting transition at 3.4 K and has strong anisotropy.  $\mu_0 H_{c2}^{\parallel ab}(T)$  and  $\mu_0 H_{c2}^{\parallel c}(T)$  curves cannot be fitted well by the WHH theory. We employed the type-II SC approximation for  $ab$ -plane and polynomial equation for  $c$ -plane. The estimated values of  $\mu_0 H_{c2}^{\parallel ab}(0)$  and  $\mu_0 H_{c2}^{\parallel c}(0)$  are 17.2 T and 0.59 T, respectively. The superconducting anisotropic parameter  $\gamma$  is determined to be 29.2.

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